

Enhancing Ferroelectrics with Insulators: Band Offsets at the $\text{Al}_2\text{O}_3/\text{BaTiO}_3$ Interface

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Engineering ferroelectric interfaces is key to improving device performance. Combining thin insulators with ferroelectrics is a promising approach for minimizing leakage currents and increasing the switching efficiency. In this work we study the interface between atomic layer deposited amorphous Al_2O_3 and epitaxial BaTiO_3 films. These materials constitute simple and robust examples of an insulator and a ferroelectric, respectively. We confirm the microstructure and interface chemistry of the heterostructure. X-ray photoelectron spectroscopy was employed to quantify the interfacial bands offsets, yielding energy barriers of 1.3 eV for holes and 2.1 eV for electrons. These results highlight Al_2O_3 as a promising candidate for an insulating layer on ferroelectrics, paving the way for efficient insulator-ferroelectrics structures for ferroelectric functional devices.

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I. INTRODUCTION

Ferroelectrics constitute some of the most studied and technologically-mature functional materials¹. In microelectronics, ferroelectrics open new routes for fast and energy-efficient devices, such as the ferroelectric field effect transistor (FeFET), a candidate for improved subthreshold performance, and a promising technology for high-performance memory devices, with applications in logic-in-memory architectures, and neuromorphic, biology-inspired computing^{2,3}. In photonics, ferroelectrics unlock efficient integrated light modulators with potential for additional active components⁴⁻⁶. In other fields ferroelectrics have shown attractive prospects in photovoltaics, catalysis, sensors and actuators⁷⁻¹¹.

A key component of the performance of functional ferroelectric devices is the ability to efficiently switch the ferroelectric layer. Two important aspects of this requirement are surface chemistry, and leakage currents through the ferroelectric. Surface chemistry can have significant influence on switching and on surface polarization¹². Leakage currents can compromise the ability to apply an electric field over the ferroelectric layer, or result in damage caused by uncontrolled currents.

Surface and interface engineering of ferroelectrics is a promising approach to overcome the above challenges. The addition of an insulating layer on the surface of a ferroelectric can help obtaining chemically well-defined interfaces, while the insulator itself can reduce leakage currents. Furthermore, adding an insulating layer on top (or the bottom) of a ferroelectric is a useful approach for improving ferroelectric tunnel junctions (FTJ) performance, where the scheme is sometimes termed ‘composite barrier’^{13,14}. The interfacial band offsets are a crucial aspect of whether an interface can be insulating¹⁵⁻¹⁷. For example, by using a capping layer of Zr doped AlO_x, the ferroelectric polarization of

HfO_{2-x} was largely improved¹⁸, and the insertion of Al₂O₃ layer into Hf_{0.5}Zr_{0.5}O₂ ferroelectrics heralded ultra-high energy storage performance¹⁹.

In recent years increasing fundamental evidence point to the crucial importance of surface chemistry and defects in determining the functional properties and performance of ferroelectrics^{20–22}. The incorporation of a surface insulating layer can change the surface chemistry in various manners, with potential effects on the ferroelectric functionality and device performance. We note that air exposure is another important facet of functional oxides^{23,24}. The current work however, focuses on the effects of an ex-situ thin Al₂O₃ layer, and as such all the BTO surfaces in this work have been exposed to the air.

Here we examine the interface between the insulator Al₂O₃ and the ferroelectric BaTiO₃ (BTO). Al₂O₃ is a common, simple and robust insulator, and BTO is one of the most studied ferroelectrics, with particular interest in integrated photonics^{25,26}. This material combination therefore combines simplicity, robustness and technological compatibility, making it a useful test case. Here we analyze the structure and interfacial chemistry and band offsets. We observe the expected surface chemistry with some non-stoichiometric Ba-rich interface phase. Importantly, we measure significant (>1 eV) band offsets at the interface, representing energy barriers for electrons and holes. These results indicate the potential of Al₂O₃ to preserve the surface chemistry of BTO and block leakage currents. Therefore, this Al₂O₃-BTO combination shows potential for improving the performance of various functional ferroelectric devices.

II. EXPERIMENTAL

8 nm of BTO were grown on TiO₂-terminated SrTiO₃ (STO) substrate (Crystal GmbH) using pulsed laser deposition (PLD) equipped with an excimer KrF laser (248nm). The laser energy and frequency were ~ 1.2 mJ/cm⁻² and 2 Hz, resulting in a growth rate of 66 pulses per monolayer. Before deposition, the substrate was heated to 670 °C using a SiC heater at a rate of 5 °C/min at 100 mTorr of oxygen, where it was annealed for 15 min prior to BTO deposition. Following deposition, the sample was annealed at the deposition temperature and oxygen pressure for 30 min. The oxygen was pumped out after the sample was cooled to 300°C. 3 nm of amorphous Al₂O₃ (alumina) layer were deposited ex-situ onto the BTO surface by atomic layer deposition (ALD) at 300 °C using trimethyl-aluminum (TMA) and H₂O (the full details are reported in a previous work¹⁵). Atomic force microscopy (AFM, Asylum MFP-3D Infinity) was used in tapping mode for topography mapping. X-ray photoelectron spectroscopy (XPS, PHI Versaprobe III) spectra were acquired using a monochromatic Al K α source (1486.6eV), and curve fitting was done using the CasaXPS software using 70:30% Gaussian/Lorentzian ratio after a Shirley-type background subtraction. Dual-beam neutralization was used to remove static charges at the surface. The binding energies were calibrated with reference to the Ti 2p 3/2 peak at 458.40 eV¹¹. X-ray diffraction (XRD, Rigaku SmartLab 9 KW) was performed using a Cu K α source ($\lambda = 1.5406$ Å) and a 2-bounce Ge monochromator, and curve fitting was done using GlobalFit 2.0 software.

III. RESULTS AND DISCUSSION

We start by characterizing the structure of the Al₂O₃/BTO bilayer, followed by its interfacial chemistry and then conclude with the band offsets. The surface topography of the Al₂O₃/BTO film measured by AFM (Figure 1) shows a step-terrace structure, with

atomically flat terraces and ~ 4 Å step height. These results indicate a high-quality growth of the BTO film as well as a smooth and uniform Al_2O_3 layer.

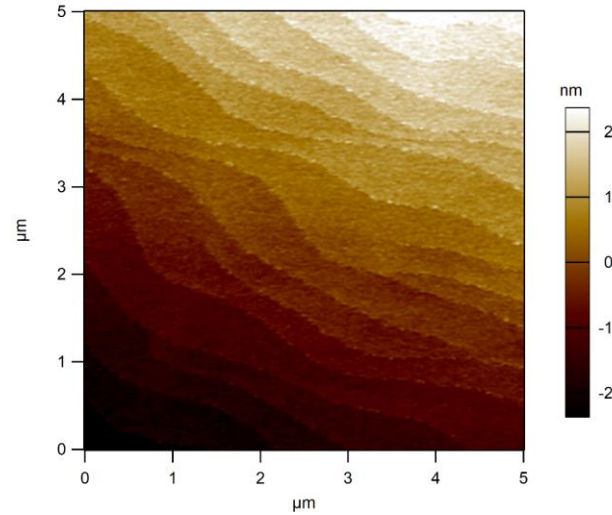


Figure 1: AFM image of the surface of 3 nm of Al_2O_3 on 8 nm of BTO

XRD analysis of the BTO film on an STO substrate is presented in Figure 2. The broad 2θ - ω scan (Figure 2a) shows the expected BTO and substrate (00L) reflections and no secondary phases, illustrating epitaxial growth of the BTO thin film on the (001)-oriented STO substrate. A feature at 38° could be the result of a miniscule amount of non-stoichiometric BTO^{27,28}, but its intensity is negligible compared to the film diffraction peaks. The ω scan of the (002) reflection (Figure 2b) exhibits a narrow FWHM of 0.17° , indicating high crystalline quality for such a thin film. Curve fitting for the out-of-plane lattice parameter yields out-of-plane lattice parameter values of 4.085 Å for the (001) peak and 4.095 Å is fitted for the (002) peak. These values are larger than the bulk of 3.994 Å (a-axis) and 4.038 Å (c-axis), suggesting that the BTO thin film is not fully strained. The lattice mismatch with the STO substrate is 2.2% for the c-oriented case, and 2.3% and 3.3% for the a-oriented case (in the two in-plane directions). In the

fully-strained BTO scenario, employing Poisson's ratio of 0.35,²⁹ the expected lattice parameters are 4.14Å (c-oriented) and 4.12Å (a-oriented), which are larger than the measured values, suggesting that the BTO is partially strained. It should be further noted that cation off-stoichiometry can cause the observed lattice parameter increase³⁰.

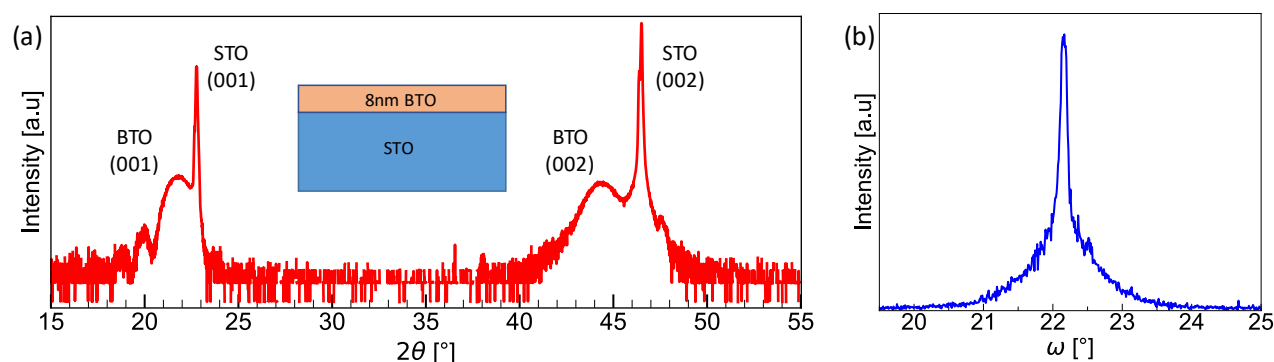


Figure 2: (a) XRD scan of the BTO/STO sample: $\theta/2\theta$, (b) ω scan of BTO/STO for the (002) peak of BTO. the inset shows a schematic drawing of the sample

XPS was acquired from the sample surface before and after deposition of 3 nm Al_2O_3 . The Ti 2p_{3/2} XPS spectra of both the BTO/STO and the Al_2O_3 /BTO surfaces (Figure 3a) correspond to Ti^{+4} , the nominal state of BTO³¹ (Table I). No under-oxidized titanium (Ti^{+3}) is observed, within the detection limit that we estimate to be better than 2% at.^{32,33} We note that for transition metals the assignment of a single formal oxidation state is not always trivial, and our discussion of Ti in this context is over-simplified for clarity³⁴. The excellent agreement between both spectra indicates that no observable redox reaction occurs during the ALD Al_2O_3 deposition process. The agreement between the Al 2p spectra of the thin (3 nm) Al_2O_3 on BTO and a thick (10 nm) Al_2O_3 from a previous work¹⁵ (Figure 3b) validates the consistency of the properties of Al_2O_3 at ultrathin thickness.

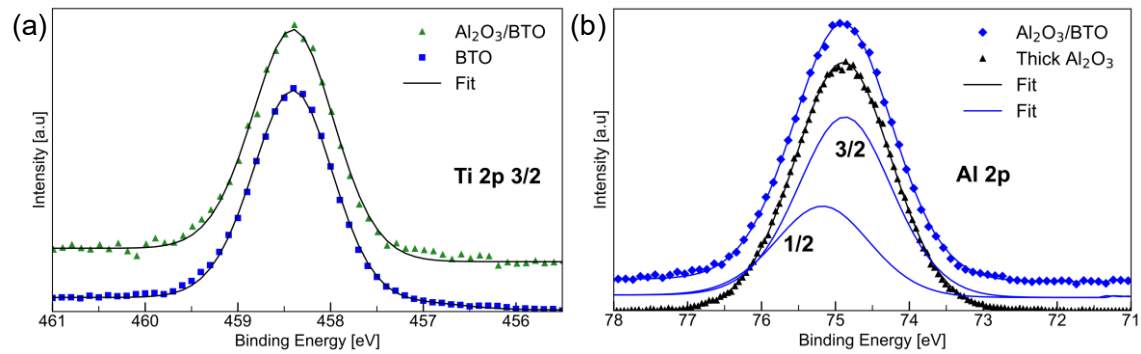


Figure 3: (a) XPS spectra of Ti 2p_{3/2} from the BTO surface, before and after ALD deposition of Al₂O₃. The binding energies are 458.40±0.05 eV. (b) XPS spectra of Al 2p from an ultrathin Al₂O₃ layer on top of a BTO, compared to a thicker Al₂O₃ from a previous work¹⁵. The binding energies of the 3/2 components are 74.52±0.05 eV and 74.77±0.05 eV, respectively. For comparison of the line shapes, the binding energy of thick Al₂O₃ was shifted by -0.25eV. The 2p 3/2 and 1/2 components of the fits are presented for Al₂O₃/BTO.

The Ba 3d doublets can each be fitted with two components. The spin-orbit value is 15.31eV in agreement with the literature³⁵. The lower-binding-energy component (Ba I) corresponds to Ba from the bulk of the BTO (Table I). The higher-binding-energy component (Ba II) is commonly attributed to surface phases or contamination of the BTO surface which may originate by BaO₂^{36,37}, indicating surface Ba excess. In both cases the Ba oxidation state is +2. The Ba II percentage of the Al₂O₃/BTO sample was lower by ~5% versus the BTO sample. Altogether the chemical state of the Al₂O₃/BTO is found to be in agreement to expectations and the literature, with confirmed well-defined single oxidation states for Ti and Al. The absence of significant foreign diffraction peaks (Figure 2) and highly smooth surface (Figure 1) indicate that this secondary phase is likely amorphous and on the BTO surface.

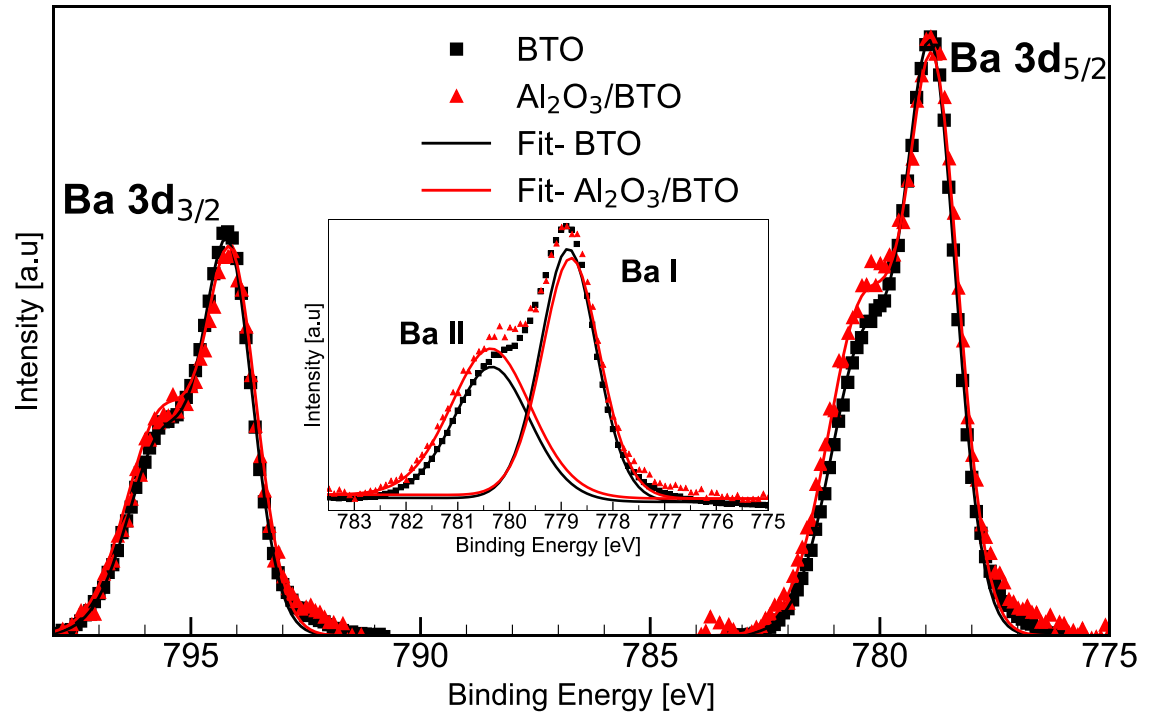


Figure 4: XPS spectra of Ba 3d of a BTO/STO and of $\text{Al}_2\text{O}_3/\text{BTO}$. The inset shows a magnified Ba $3d_{5/2}$ region with the fits of Ba I and Ba II components (see text for details).

Table I. The binding energy values of the core levels before and after Al_2O_3 deposition. The values have an uncertainty level of ± 0.05 eV

Sample	Al 2p _{3/2} [eV]	Ti 3p [eV]	Ti 2p _{3/2} [eV]	Ba 3d 5/2 [eV]		Ba 3d 3/2 [eV]	
				Ba I	Ba II	Ba I	Ba II
bare BTO	-	36.28	458.40	778.86	780.34	794.17	795.65
$\text{Al}_2\text{O}_3/\text{BTO}$	74.52	36.37	458.40	778.81	780.36	794.10	795.65

Having ascertained the structure and interface chemistry of the $\text{Al}_2\text{O}_3/\text{BTO}$ structure, we turn the focus to the interfacial band offsets. The valence band offset (VBO) at the $\text{Al}_2\text{O}_3/\text{BTO}$ interface was determined using the Kraut method^{38,39}

$$(1) \text{VBO} = \text{VBE}_{\text{BTO}} - \text{VBE}_{\text{Al}_2\text{O}_3} = A - B - C$$

$$(2) A = \text{BE}_{\text{Ti } 2p}^{\text{BTO}} - \text{VBE}_{\text{BTO}} = 455.9 \text{ eV}$$

$$(3) B = \text{BE}_{\text{Ti } 2p}^{\text{Al}_2\text{O}_3/\text{BTO}} - \text{BE}_{\text{Al } 2p \ 3/2}^{\text{Al}_2\text{O}_3/\text{BTO}} = 383.9 \text{ eV}$$

$$(4) C = BE_{Al\ 2p\ 3/2}^{Al_2O_3} - VBE_{Al_2O_3} = 70.7\text{eV}$$

Where VBE_X is the valence band edge of material X (BTO or Al_2O_3), respectively, and BE_X^Y is the binding energy of the element X in material Y. The analysis yields a VBO of 1.3 ± 0.2 eV. Using the known values of the band gaps for BTO and Al_2O_3 the conduction band offset (CBO) was then calculated as:

$$(5) CBE_{Al_2O_3} - CBE_{BTO} = E_g^{Al_2O_3} - E_g^{BTO} - (VBE_{BTO} - VBE_{Al_2O_3})$$

Where CBE_X is the conduction band edge of material X, E_g^X is the band gap of material X. The band gap value of Al_2O_3 used in this work (6.6 eV) corresponds to Al_2O_3 grown under comparable conditions¹⁵. The resulting CBO is 2.1 ± 0.3 eV^{37–39}. The calculation was repeated using the Ti 3p (instead of Ti 2p 3/2), and it yielded consistent results within the margin of error. These values indicate that Al_2O_3 is a suitable insulating layer on top of BTO from a band structure perspective. The band structure (Figure 6) illustrates the interfacial barrier for both electrons and holes. The band gap value used for BTO corresponds to its bulk form and may differ when considering a thin film⁴².

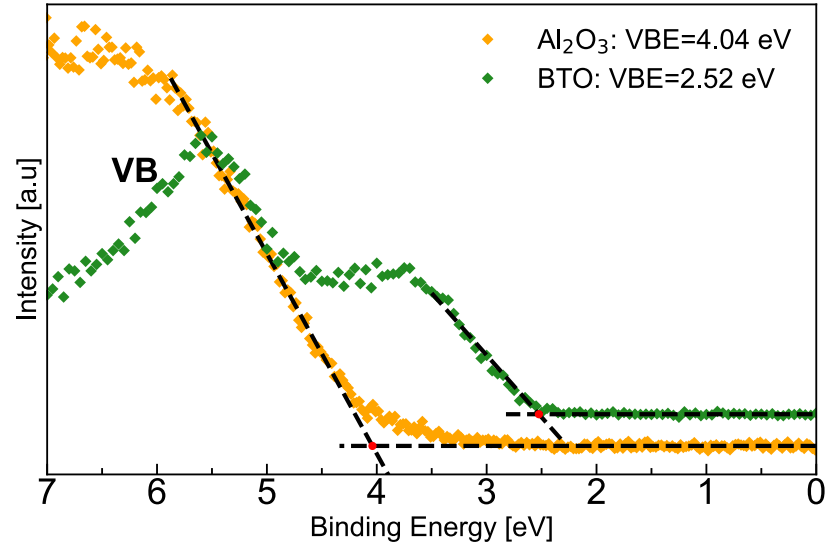


Figure 5: XPS spectrum of the valence band edges of the BTO and of Al_2O_3

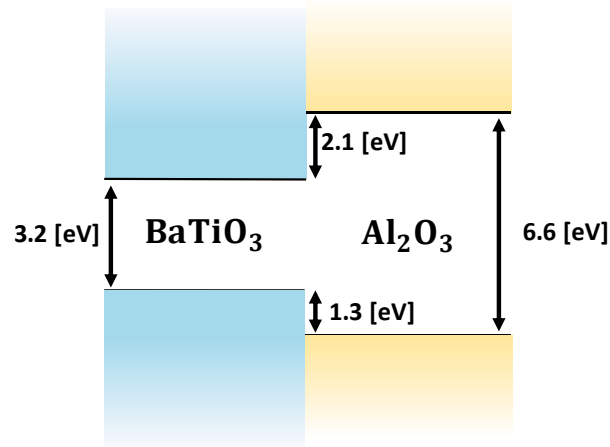


Figure 6: The band structure of BTO- Al_2O_3 interface

IV. SUMMARY AND CONCLUSIONS

The interface between Al_2O_3 and BTO was investigated using XPS. The analysis reveals that ex situ ALD growth of amorphous Al_2O_3 on BTO results in an insulating layer, with interfacial barriers of 1.3 eV for holes and 2.1 eV for electrons. The deposition process preserves the surface chemistry of BTO. These results suggest that Al_2O_3 is a promising insulating material for integration with ferroelectric BTO in electronic devices.

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AUTHOR DECLARATIONS

Conflicts of Interest

The authors have no conflicts to disclose.

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